PERFORMANCE SIMULATION OF A NATURAL CIRCULATION SOLAR AIR HEATER WITH PHASE CHANGE MATERIAL ENERGY STORAGE

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Abstract
The one-dimensional radial heat conduction in a phase change material (PCM) heat storage system encapsulated in cylindrical pipes in a single glazed flat plate natural circulation solar air heater is presented. The PCM is prepared in modules, with the modules equispaced across the absorber plate. Enthalpy method is applied to convert the relevant energy balance equations into dimensionless forms for easy tracking of the moving phase boundaries. Crank-Nicolson implicit finite difference scheme which has the features of being stable, accurate and fast in its solution is used in the solution of the governing equations subject to suitable initial and boundary conditions. The scheme is applied at each node and the resulting simultaneous equations are solved using the Gauss-Seidel iterative method. An existing computer programme in BASIC known as the EGGINC which was developed for the rectangular channel containing the PCM is modified to cylindrical coordinate for the pipes containing the PCM to predict the temperature distributions in the solar air heater. The predicted temperatures of the system is compared with the experimental data under daytime no-load condition over the ambient temperature range of 18.5-36.0^{O}C and daily global irradiation of 4.9-20.1MJ/m^2-day. The predicted temperatures agree closely with experimental data to within acceptable limits.

Keywords: Phase change material(PCM); Natural circulation; Solar air heater; Enthalpy method; Dimensionless forms; Crank-Nicolson implicit finite difference scheme; Predicted temperatures; Experimental data

1.0 Introduction

   The heat-conduction equation and its solutions express temperature as a function of the space coordinate and time. This work deals with a survey of existing numerical technique – the Crank-Nicolson implicit finite difference scheme for solving one-dimensional radial problem of a tube containing phase change material in a natural circulation solar air heater used for many purposes. In particular we will consider the melting/freezing problem using the enthalpy method, hence, the temperature range in the phase change using the heat conduction equation. Crank-Nicolson method provides a good introduction in the phase change problem and presents an elaborate approximation for these problems. The implicit methods are the natural alternative to the front tracking methods. Within these implicit methods the most widely used are the enthalpy method. In the enthalpy method described in Voller and Cross (1981), the total enthalpy of the system, is utilized.

   In many industrial problems the phase change occurs over a temperature range rather than at a specified temperature. The enthalpy method can also be applied to such problems Enibe (2003).These methods appear to have great flexibility and are easily extended to multidimensional problems. The advantages of this approach are stated in (Voller and Cross1981). The solar air heater is considered to be a promising direction for increasing the economic feasibility of low temperature solar systems for heating water for domestic, agricultural and industrial applications. A system of this type combines collection and storage of thermal energy in a single unit (Rabin et al, 1995).Most investigators have devoted greater attention to forced circulation air heaters operating under steady state conditions. In contrast, only few reports of natural circulation air heaters have appeared in the literature (Macedo and Altemani, 1978).

   In this work, cylindrical tubes or pipes encapsulating phase change materials (wax) are used instead of rectangular blocks or channels reported in Enibe (2003). In the previous work the analysis was purely on rectangular coordinates whereas herein cylindrical coordinate is used extensively in the analysis. This work will be undertaken using an upgraded computer simulation program, based on an optimally verified transient thermal analysis to validate and optimise the system. Natural circulation air heaters are important in many industrial and agricultural applications including the drying of crops and medicinal/aromatic plants, timber, natural rubber, tea and coffee products, and fodder for animals (Ekechukwu and Norton, 1998; Diamante and Munro, 1993; Fohr and Figueiredo, 1980; Palaniappan and Subramaian, 1998). It could also be used for poultry egg incubation. In the latter case, as well as in the drying of medicinal/aromatic plants, the heated air temperatures is to be maintained within specified ranges. Further, the hot air is to be supplied over a continuous period of several days, including off-sunshine periods. For these special applications, some form of energy

storage, possibly combined with an auxiliary heat source is required. Other applications of the natural circulation solar air heaters are in domestic hot water and space heating, and to a lesser degree, in industrial processes. They can be employed to supply hot water in absorption refrigeration systems for space cooling.

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2.0 Derivation of the Simultaneous Equations for the Temperature Distributions in a Solar Air Heater

The 1 – D differential equation governing heat transfer in cylindrical coordinates is given by

$$\rho_s c_s \frac{\partial T_s}{\partial t} = k_s \left\{ \frac{1}{r} \frac{\partial T_s}{\partial r} + \frac{\partial^2 T_s}{\partial r^2} \right\} \tag{2.1}$$

We assume that the heat flows in the y and z directions are negligible.

The initial condition is

$$T_s (r,0) = f(r), 0 \leq r \leq r_m \tag{2.2}$$

Where $f(r) = T_s = T_{ambient}$

The boundary conditions are as follows:

(i)   At the edge of the pipe, the radius $r = r_m$ is maximum and there is continuation of heat flux.

$$\therefore - k_s \frac{\partial T_s}{\partial r} = U_s \left(T_{p1} - T_s \right) \tag{2.3}$$

(ii)   At the centre of the pipe, the radius is $r = 0$ and there is temperature symmetry.

$$\therefore \frac{\partial T_s}{\partial r} = 0 \tag{2.4}$$

To overcome the problems of tracking the moving boundary associated with the melting/freezing of the phase change material, we employ the total enthalpy method described in Voller and Cross (1981) and Enibe (2003). Here, the heat transfer equation is rewritten in terms of total enthalpy of the phase change material as

$$\rho_s \frac{\partial H_s}{\partial t} = \frac{1}{r} \frac{\partial T_s}{\partial r} + \frac{\partial^2 T_s}{\partial r^2} \tag{2.5}$$

The phase change material used does not melt at a single temperature, but rather over a temperature range say T_1 and T_2. The temperature – enthalpy relationships for the material may be represented as a series of straight lines as shown in table 1 and fig.1.

Table 1.: Temperature-Enthalpy Relationship For Paraffin Wax(PCM)

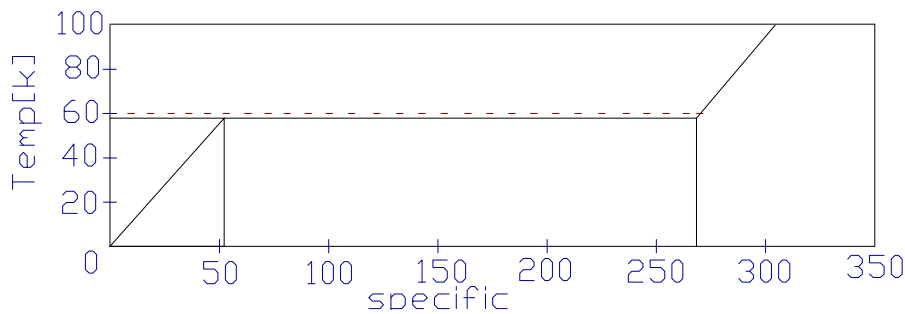| Temperature (K)	Specific	Enthalpy (KJ kg^{-1})
273	0	
331	52.2	H_1
333	268.4	H_2
373	304.4	

Fig 1. Temperature-Enthalpy Relationship for Paraffin Wax(PCM)

Suppose the specific enthalpy of the PCM at a reference temperature T_o is H_o, the specific enthalpy of the fully solid phase at a minimum melting temperature T_1 is H_1, while the specific enthalpy of the liquid phase at the maximum melting temperature T_2 is H_2. The specific enthalpies are calculated from the expressions presented in Enibe (2003) as follows:

$$H_1 = c_s \left(T_1 - T_0 \right) \tag{2.6}$$

$$H_2 = H_1 + c_s \left(T_2 - T_1 \right) + \Delta H_s' = H_1 + \Delta H_s \tag{2.7}$$

where T_0, is an arbitrary reference temperature, and $\Delta H_s = \Delta H_s' + c_s \left(T_2 - T_1 \right)$, the modified enthalpy of fusion. We assume that at T_1, the PCM is fully solid, while at T_2, it is fully liquid. In between the two temperatures, the liquid fraction and the temperature vary linearly with specific enthalpy.

For any temperature less than T_1, the specific enthalpy is given by the expression

$$H_s = H_0 + c_s \left(T_s - T_0 \right), \quad T_s < T_1 \tag{2.8}$$

Suppose the liquid fraction at any point between the minimum and maximum melting temperature is α. Then, the specific enthalpy of a mixture of the liquid and solid phase at any temperature, T_s is given by

$$H_s = H_1 + \alpha \Delta H_s \tag{2.9}$$

$$T_s = T_1 + \alpha \Delta T_s = T_1 + \alpha \left(T_2 - T_1 \right), \quad T_1 \leq T_s \leq T_2 \tag{2.10}$$

Above the maximum melting temperature T_2, the PCM is fully liquid and the specific enthalpy is given by

$$H_s = H_2 + c_s \left(T_s - T_2 \right), \quad T_s > T_2 \tag{2.11}$$

Since T_0 can be taken quite arbitrary, we set $H_0 = 0$ at $T_0 = 273k$.

To simplify the utilization of the temperature-enthalpy relationships in the governing equations, we define the dimensionless enthalpy, Ψ_s as

$$\Psi_s = \frac{H_s - H_0}{\Delta H_s} = \frac{H_s}{\Delta H_s} \tag{2.12}$$

Let $\Psi_1 = \dfrac{H_1}{\Delta H_s}$ \hfill (2.13)

and

$$\Psi_2 = \frac{H_2}{\Delta H_s} \tag{2.14}$$

For the range $T_s < T_1$, and from equations (2.8) and (2.12) we have

$$\therefore T_s = \frac{\psi_s \Delta H_s}{c_s} + T_0 \tag{2.15}$$

For the range $T_1 \leq T_s \leq T_2$ we have from equation (2.9) that

$$H_s = H_1 + \alpha \Delta H_s, \quad T_1 \leq T_s \leq T_2$$

But $\alpha = \dfrac{T_s - T_1}{T_2 - T_1} = \dfrac{T_s - T_1}{2 \in_s}$ \hfill (2.16)

Hence, $H_s = H_1 + \dfrac{T_s - T_1}{2 \in_s} \Delta H_s$ \hfill (2.17)

Dividing through by ΔH_s, we have

$$\Psi_s = \Psi_1 + \frac{T_s - T_1}{2 \in_s} \tag{2.18}$$

$$\Rightarrow 2 \in_s \left(\Psi_s - \Psi_1 \right) = T_s - T_1 \tag{2.19}$$

For the range $T_s > T_2$, we have from equation (2.11) that

$$H_s \;=\; H_2 \;+\; c_s \left(T_s - T_2\right), \quad T_s > T_2$$

Since $H_s \;=\; H_0 \;+\; \Psi_s \Delta H_s$

Where H$_0$ is set equal to zero, and $H_2 \;=\; \Psi_2 \Delta H_s$, we obtain

$$T_s = \frac{\Delta H_s}{c_s}\left(\psi_s - \psi_2\right) + T_2 \tag{2.20}$$

From equations (2.15), (2.19) and (2.20) we can write in general that

$$T \;=\; \lambda_1 \psi \;+\; \lambda_2 \tag{2.21}$$

Where λ_1, λ_2 are constants which depends on the temperature range.

From equation (2.22), we can write more specifically for the phase change material as

$$T_s \;=\; \lambda_{1s} \psi_s \;+\; \lambda_{2s} \tag{2.22}$$

Where it is understood that λ_{1s} and λ_{2s} are evaluated at T$_s$. From equation (2.21) and equation (2.22), the difference between the temperature of any material and that of the phase change material, T – T$_s$ may be given by

$$T \;-\; T_s \;=\; \lambda_1 \psi \;+\; \lambda_2 \;-\; \lambda_{1s} \psi_s \;-\; \lambda_{2s}$$

$$T \;-\; T_s \;=\; \lambda_2 \left[\psi + \left(\lambda_2 - \lambda_{1s} \psi_s - \lambda_{2s}\right)\big/\lambda_1\right] \tag{2.23}$$

$$\therefore\; T \;-\; T_s \;=\; \lambda_1 \left(\psi - \psi_s^*\right) \tag{2.24}$$

Where $\psi_s^* \;=\; \left(\psi_s \lambda_{1s} + \lambda_{2s} - \lambda_2\right)\big/\lambda_1$ \hfill (2.25)

On the other hand, for any two materials not involving the phase charge material such as T$_{f1}$ and T$_{f2}$, λ_1 and λ_2 are constants and equal, hence, the temperature difference is given by

$$T_{f1} \;-\; T_{f2} \;=\; \lambda_1 \left(\psi_{f1} - \psi_{f2}\right) \tag{2.26}$$

To simplify the numerical solution of the equation and also aid in the parametric analysis of the natural circulation, solar air heater, we define the following dimensionless parameters.

$$R \;=\; \frac{r}{r_m} \tag{2.27}$$

$$\tau \;=\; \frac{t}{t_o} \tag{2.28}$$

Where t$_0$ is some arbitrary time interval

Utilizing the above dimensionless enthalpy terms, the differential coefficients may be given as follows:

$$\frac{\partial H_s}{\partial t} \;=\; \frac{\Delta H_s}{t_0}\;\frac{\partial \psi_s}{\partial \tau} \tag{2.29}$$

$$\frac{\partial T_s}{\partial t} \;=\; \frac{\lambda_1}{t_0}\;\frac{\partial \psi_s}{\partial \tau} \tag{2.30}$$

$$\frac{1}{r}\;\frac{\partial T_s}{\partial r} \;=\; \frac{\lambda_1}{r_m^2}\;\frac{1}{R}\;\frac{\partial \psi_s}{\partial R} \tag{2.31}$$

$$\frac{\partial T_s}{\partial r} \;=\; \frac{\lambda_1}{r_m}\;\frac{\partial \psi_s}{\partial R} \tag{2.32}$$

$$\frac{\partial^2 T_s}{\partial r^2} \;=\; \frac{\lambda_1}{r_m^2}\;\frac{\partial^2 \psi_s}{\partial R^2} \tag{2.33}$$

Where r$_m$ is the inside radius of the pipe containing the PCM. It is noted that for the phase change material, λ_1 and λ_2 are replaced by λ_{1s} and λ_{2s}.

Substituting the dimensionless coefficients into the governing equation for the phase change material, simplifying and multiplying through by $r_m^2\big/\lambda_{1s}$ to obtain

$$\frac{\rho_s}{k_s} \; \frac{\Delta H_s r_m^2}{\lambda_{1s} t_0} \; \frac{\partial \psi_s}{\partial \tau} \;\; = \;\; \frac{1}{R} \; \frac{\partial \psi_s}{\partial R} \;\; + \;\; \frac{\partial^2 \psi_s}{\partial R^2} \tag{2.34}$$

Now let

$$\lambda_{1s} \;\; = \;\; \lambda_1 \; \frac{\lambda_{1s}}{\lambda_1} \;\; = \;\; \lambda_1 \; \beta_s \tag{2.35}$$

Using the expressions for λ_1 and λ_{1s} from the temperature - enthalpy relationship, we have

$$\beta_s \;\; = \;\; \frac{\Delta H_s}{c_s} \; \frac{c_s}{\Delta H_s} \;\; = \;\; 1, \quad\quad T_s \;\; < \;\; T_1 \tag{2.36}$$

$$\beta_s \;\; = \;\; \frac{2 \,\epsilon_s}{1} \; \frac{c_s}{\Delta H_s} \;\; = \;\; \frac{2 \,\epsilon_s \, c_s}{\Delta H_s}, \quad\quad T_1 \;\; \leq \;\; T_s \;\; \leq \;\; T_2 \tag{2.37}$$

$$\beta_s \;\; = \;\; \frac{\Delta H_s}{c_s} \; \frac{c_s}{\Delta H_s} \;\; = \;\; 1, \quad\quad T_s \;\; > \;\; T_2 \tag{2.38}$$

It is noted that for a phase change material with constant phase change temperature, $T_1 \;\; = \;\; T_2 \;\; = \;\; T_m$, $\epsilon_s \;\; = \;\; \left(T_2 \, - \, T_1\right)\big/2 \;\; = \;\; 0$ and hence $\beta_s \;\; = \;\; 0$. Thus, β_s is a phase change temperature range factor. Since t_0 can be taken quite arbitrary, we set

$$\frac{\rho_s}{k_s} \; \frac{\Delta H_s r_m^2}{\lambda_{1s} t_0} \;\; = \;\; \frac{\rho_s \Delta H_s \, r_m^2 \, c_s}{k_s \, \Delta H_s \, t_0} \;\; = \;\; \frac{\rho_s \, c_s \, r_m^2}{k_s \, t_0} \;\; = \;\; 1 \;\text{ to get } t_0 \;\; = \;\; \frac{r_m^2 \rho_s \, c_s}{k_s} \tag{2.39}$$

Thus, $t_0 \;\; = \;\; \dfrac{r_m^2 \rho_s \, c_s}{k_s} \tag{2.40}$

Substituting t_0 and λ_{1s} into equation (2.34) and simplifying to obtain

$$\frac{\rho_s \Delta H_s \, r_m^2 k_s \, c_s}{k_s \, \beta_s \, \Delta H \rho_s \, r_m^2 \, c_s} \; \frac{\partial \psi_s}{\partial \tau} \;\; = \;\; \frac{1}{R} \; \frac{\partial \psi_s}{\partial R} \;\; + \;\; \frac{\partial^2 \psi_s}{\partial R^2}$$

To give

$$\frac{\partial \psi_s}{\partial \tau} \;\; = \;\; \beta_s \; \left\{ \frac{1}{R} \; \frac{\partial \psi_s}{\partial R} \;\; + \;\; \frac{\partial^2 \psi_s}{\partial R^2} \right\} \tag{2.41}$$

The boundary conditions then become

(i) $\quad\quad r \;\; = \;\; r_m, \quad R \;\; = \;\; 1$

We have

$$- \; k_s \; \frac{\partial T_s}{\partial r} \;\; = \;\; u_s \; \left(T_{p1} \, - \, T_s\right) \text{ equation (2.3)}$$

but $- \; k_s \; \dfrac{\partial T_s}{\partial r} \;\; = \;\; - \; k_s \; \dfrac{\lambda_{1s}}{r_m} \; \dfrac{\partial \psi_s}{\partial R} \tag{2.42}$

$$U_s \; T_{p1} \;\; = \;\; U_s \, \lambda_1 \; \psi_{p_1} \tag{2.43}$$

$$U_s \; T_s \;\; = \;\; U_s \, \lambda_1 \; \psi_s \tag{2.44}$$

$$\therefore \;\; - \; k_s \; \frac{\lambda_{1s}}{r_m} \; \frac{\partial \psi_s}{\partial R} \;\; = \;\; U_s \, \lambda_1 \; \left(\psi_{p1} \, - \, \psi_s\right) \tag{2.45}$$

$$\therefore \;\; - \; k_s \; \frac{\lambda_{1s}}{r_m} \; \frac{\partial \psi_s}{\partial R} \;\; = \;\; U_s \, \lambda_1 \; f\left(\psi_p \,, \, \psi_s\right) \tag{2.46}$$

Where $f\left(\psi_{p1} \,, \, \psi_s\right)$ = functions of dimensionless enthalpy of plate 1 or temperature of plate 1 and dimensionless enthalpy of phase change material or temperature of phase change material.

Simplifying to obtain

$$\frac{\partial \psi_s}{\partial R} = -\frac{U_s\, r_m}{k_s} \frac{\lambda_1}{\lambda_{1s}} f\left(\psi_{p1}\, ,\ \psi_s\right)$$

$$\frac{\partial \psi_s}{\partial R} = -\frac{U_s\, r_m}{k_s} \frac{1}{\beta_s} f\left(\psi_{p1}\, ,\ \psi_s\right)$$

$$\therefore \quad \frac{\partial \psi_s}{\partial R} = -\frac{S'_{r_1}}{\beta_s} \tag{2.47}$$

Discretizing equation (2.47) we have

$$\frac{1}{4(\Delta R)}\left\{ \begin{array}{l} \psi_s\left(i+1,\, j+1\right) + \psi_s\left(i+1,\, j\right) \\ -\ \psi_s\left(i-1,\, j+1\right) - \psi_s\left(i-1,\, j\right) \end{array} \right\} = -\frac{S'_{r_m}}{\beta_s} \tag{2.48}$$

Since i = m, we have

$$\psi_s\left(m+1,\, j+1\right) + \psi_s\left(m+1,\, j\right)$$

$$= \psi_s\left(m-1,\, j+1\right) + \psi_s\left(m-1,\, j\right) - \frac{4(\Delta R)\, S'_{r_m}}{\beta_s} \tag{2.49}$$

where

ψ_{p1} = Plate dimensionless enthalpy and

ψ_s = Phase change dimensionless enthalpy

(ii)   Similarly, at r = 0, R = 0 and $\dfrac{\partial \psi}{\partial R} = 0$

$$\frac{\partial \psi_s}{\partial R} = 0 \tag{2.50}$$

Discretizing equation (2.50) we obtain

$$\frac{1}{4(\Delta R)}\left\{ \begin{array}{l} \psi_s\left(i+1,\, j+1\right) + \psi_s\left(i+1,\, j\right) \\ -\ \psi_s\left(i-1,\, j+1\right) - \psi_s\left(i-1,\, j\right) \end{array} \right\} = 0 \tag{2.51}$$

Since i = 0, we have

$$\psi_s\left(1,\, j+1\right) + \psi_s\left(1,\, j\right) = \psi_s\left(-1,\, j+1\right) + \psi_s\left(-1,\, j\right) \tag{2.52}$$

In order to achieve stable, accurate and fast solutions of the governing equations subject to the specified boundary conditions, Crank-Nicolson implicit finite difference scheme is utilized. The scheme, describe in Kreyszig (2006) and Carnaham et al (1969), has the advantage of being unconditionally stable, and has an accuracy of second order compared with the explicit method. The technique involves sweeping across time and space interval in one step. The resulting algebraic equations will therefore contain more than three unknowns at any given node in a time step. Consequently, the tridiagonal algorithm cannot be used alone to solve the equations. We therefore employ the Gauss-Seidel iterative scheme to solve the resulting equation in the tridiagonal algorithm using iterations.

For the numerical solution the dimensionless forms of the partial differential coefficient are as follows:

$$\frac{\partial \psi_s}{\partial \tau} = \frac{1}{\Delta \tau} \left\{ \psi_s\left(i,\, j+1\right) - \psi_s\left(i,j\right) \right\} \tag{2.53}$$

$$R = i\left(\Delta R\right) \tag{2.54}$$

$$\frac{\partial \psi_s}{\partial R} = \frac{1}{4(\Delta R)}\left\{ \begin{array}{l} \psi_s\left(i+1,\, j+1\right) + \psi_s \\ \left(i+1,\, j\right) - \psi_s\left(i-1,\, j+1\right) - \psi_s\left(i-1,\, j\right) \end{array} \right\} \tag{2.55}$$

$$\frac{\partial^2 \psi_s}{\partial R^2} = \frac{1}{2(\Delta R)^2} \left\{ \begin{array}{l} \psi_s\left(i+1, j+1\right) + \psi_s\left(i+1, j\right) - 2\psi_s\left(i, j+1\right) \\ -2\psi_s\left(i, j\right) + \psi_s\left(i-1, j+1\right) + \psi_s\left(i-1, j\right) \end{array} \right\} \tag{2.56}$$

Substituting the finite difference form of the differential coefficient into the governing equation, we have that for any given pipe in the solar air heater, the finite difference forms of the equation (from equations 2.41, 2.53, 2.54, 2.55, and 2.56) become

$$\frac{\beta_s}{i(\Delta R)} \cdot \frac{1}{4(\Delta R)} \left\{ \begin{array}{l} \psi_s\left(i+1, j+1\right) + \psi_s\left(i+1, j\right) \\ -\psi_s\left(i-1, j+1\right) - \psi_s\left(i-1, j\right) \end{array} \right\} +$$

$$\frac{\beta_s}{2(\Delta R)^2} \left\{ \begin{array}{l} \psi_s\left(i+1, j+1\right) + \psi_s\left(i+1, j\right) - 2\psi_s\left(i, j+1\right) \\ -2\psi_s\left(i, j\right) + \psi_s\left(i-1, j+1\right) + \psi_s\left(i-1, j\right) \end{array} \right\} =$$

$$\frac{1}{\Delta \tau} \left\{ \psi_s\left(i, j+1\right) - \psi_s\left(i, j\right) \right\} \tag{2.57}$$

Collecting like terms and simplifying, become

$$\psi_s\left(i+1, j+1\right) \frac{\beta_s}{2(\Delta R)^2} \left\{ \frac{1}{2i} + 1 \right\} - \psi_s\left(i, j+1\right)$$

$$\left\{ \frac{1}{\Delta \tau} + \frac{\beta_s}{(\Delta R)^2} \right\} + \psi_s\left(i-1, j+1\right)$$

$$\frac{\beta_s}{2(\Delta R)^2} \left\{ 1 - \frac{1}{2i} \right\} = -\psi_s\left(i+1, j\right) \frac{\beta_s}{2(\Delta R)^2}$$

$$\left\{ 1 + \frac{1}{2i} \right\} + \psi_s\left(i, j\right) \left\{ \frac{\beta_s}{(\Delta R)^2} - \frac{1}{\Delta \tau} \right\}$$

$$+ \psi_s\left(i-1, j\right) \frac{\beta_s}{2(\Delta R)^2} \left\{ \frac{1}{2i} - 1 \right\} \tag{2.58}$$

This may be written compactly in the tridiagonal format as

$$A_1\left(i, j\right) \psi_s\left(i+1, j+1\right) + B_1\left(i, j\right)$$

$$\psi_s\left(i, j+1\right) + C_1\left(i, j\right) \psi_s\left(i-1, j+1\right) = D_1\left(i, j\right) \tag{2.59}$$

Where A_1, B_1, and C_1 has their significance.
The boundary conditions are
(i) At the outermost part of the pipe $r = r_m$, $i = m$; the finite difference formulation of the boundary condition from equation (2.49) is

$$\frac{1}{4(\Delta R)} \left\{ \begin{array}{l} \psi_s\left(i+1, j+1\right) + \psi_s\left(i+1, j\right) \\ -\psi_s\left(i-1, j+1\right) - \psi_s\left(i-1, j\right) \end{array} \right\} = -\frac{S'_{r_m}}{\beta_s}$$

Let $r = r_m$, $i = m$, we have

$$\frac{\beta_s}{4(\Delta R)} \left\{ \begin{array}{l} \psi_s\left(m+1, j+1\right) + \psi_s\left(m+1, j\right) \\ -\psi_s\left(m-1, j+1\right) - \psi_s\left(m-1, j\right) \end{array} \right\} = -S'_{r_m} \tag{2.60}$$

From equation (2.52) we have

$$\psi_s\left(m+1, j+1\right) + \psi_s\left(m+1, j\right) = \psi_s\left(m-1, j+1\right) + \psi_s\left(m-1, j\right) - \frac{4(\Delta R) S'_{r_m}}{\beta_s}$$

Using the general node equation, equation (2.61), substituting $i = m$ and simplifying further we have

$$\left[\psi_s\left(m+1,\,j+1\right)+\psi_s\left(m+1,\,j\right)\right]\frac{\beta_s}{2(\Delta R)^2}\left\{1+\frac{1}{2m}\right\}-\psi_s\left(m,\,j+1\right)$$

$$\left\{\frac{1}{\Delta\tau}+\frac{\beta_s}{(\Delta R)^2}\right\}+\psi_s\left(m-1,\,j+1\right)$$

$$\frac{\beta_s}{(\Delta R)^2}\left\{1-\frac{1}{2m}\right\}=\psi_s\left(m,j\right)\left\{\frac{\beta_s}{(\Delta R)^2}-\frac{1}{\Delta\tau}\right\}$$

$$+\psi_s\left(m-1,\,j\right)\frac{\beta_s}{2(\Delta R)^2}\left\{\frac{1}{2m}-1\right\}$$

(2.61)

Substituting equation (2.49) into equation (2.61) and after rearranging to get

$$\psi_s\left(m-1,\,j+1\right)\frac{\beta_s}{2(\Delta R)^2}\left\{\left(1+\frac{1}{2m}\right)+\left(1-\frac{1}{2m}\right)\right\}$$

$$+\psi_s\left(m-1,\,j\right)\frac{\beta_s}{2(\Delta R)^2}\left\{\left(1+\frac{1}{2m}\right)-\left(1-\frac{1}{2m}\right)\right\}$$

$$-\psi_s\left(m,\,j+1\right)\left\{\frac{1}{\Delta\tau}+\frac{\beta_s}{(\Delta R)^2}\right\}=\psi_s\left(m,\,j\right)\left\{\frac{\beta_s}{(\Delta R)^2}-\frac{1}{\Delta\tau}\right\}+\frac{2S'_{r_m}}{(\Delta R)}$$

(2.62)

Solving and simplifying we have

$$\psi_s\left(m-1,\,j+1\right)\frac{\beta_s}{(\Delta R)^2}-\psi_s\left(m,\,j+1\right)\left\{\frac{1}{\Delta\tau}+\frac{\beta_s}{(\Delta R)^2}\right\}=$$

$$-\psi_s\left(m-1,\,j\right)\frac{\beta_s}{2m(\Delta R)^2}+\psi_s\left(m,\,j\right)\left\{\frac{\beta_s}{(\Delta R)^2}-\frac{1}{\Delta\tau}\right\}+\frac{2S'_{r_m}}{\Delta R}$$

(2.63)

In tridiagonal terms, this may be rewritten as

$$A\left(m,\,j\right)\psi_s\left(m-1,\,j+1\right)+B(m,j)\psi_s\left(m,\,j+1\right)=D(m,j)$$

(2.64)

Where A(m , j),B(m , j),and D (m , j) have their significance

Similarly, at the centre of the pipe, r = 0, the finite difference formulation of the boundary condition from equation (2.51) becomes

$$\frac{1}{4(\Delta R)}\left\{\begin{array}{l}\psi_s\left(i+1,\,j+1\right)+\psi_s\left(i+1,\,j\right)\\-\psi_s\left(i-1,\,j+1\right)-\psi_s\left(i-1,\,j\right)\end{array}\right\}=0$$

$$\therefore\psi_s\left(i+1,\,j+1\right)+\psi_s\left(i+1,\,j\right)$$

$$-\psi_s\left(i-1,\,j+1\right)+\psi_s\left(i-1,\,j\right)=0$$

(2.65)

For $r=0$ $\therefore i=0$

Hence

$$\psi_s\left(i,\,j+1\right)+\psi_s\left(1,\,j\right)-\psi_s\left(-1,\,j+1\right)-\psi_s\left(-1,\,j\right)=0$$

(2.66)

From equation (2.52), we have

$$\psi_s\left(1,\,j+1\right)+\psi_s\left(1,\,j\right)=\psi_s\left(-1,\,j+1\right)+\psi_s\left(-1,\,j\right)$$

$$\Rightarrow\psi_s\left(-1,\,j+1\right)+\psi_s\left(-1,\,j\right)=\psi_s\left(1,\,j+1\right)+\psi_s\left(1,\,j\right)$$

(2.67)

From equation (2.58) the general node equation, and substituting i = 0, we obtain

$$\psi_s\left(1, j+1\right)\frac{\beta_s}{2(\Delta R)^2} - \psi_s\left(0, j+1\right)\left\{\frac{1}{\Delta\tau} + \frac{\beta_s}{(\Delta R)^2}\right\} + \psi_s\left(-1, j+1\right)$$

$$\frac{\beta_s}{2(\Delta R)^2} = -\psi_s\left(1, j\right)\frac{\beta_s}{2(\Delta R)^2}$$

$$+ \psi_s\left(0, j\right)\left\{\frac{\beta_s}{(\Delta R)^2} - \frac{1}{\Delta\tau}\right\} - \psi_s\left(-1, j\right)\frac{\beta_s}{2(\Delta R)^2} \quad (2.68)$$

Rearranging and substituting equation (2.67) into equation (2.68), we have

$$\psi_s\left(1, j+1\right)\frac{\beta_s}{2(\Delta R)^2} - \psi_s\left(0, j+1\right)$$

$$\left\{\frac{1}{\Delta\tau} + \frac{\beta_s}{(\Delta R)^2}\right\} + \left\{\psi_s\left(1, j+1\right) + \psi_s\left(1, j\right)\right\}\frac{\beta_s}{2(\Delta R)^2} =$$

$$-\psi_s\left(1, j\right)\frac{\beta_s}{2(\Delta R)^2} + \psi_s\left(0, j\right)\left\{\frac{\beta_s}{(\Delta R)^2} - \frac{1}{\Delta\tau}\right\} \quad (2.69)$$

Solving and simplifying, we have

$$\psi_s\left(1, j+1\right)\left\{\frac{\beta_s}{2(\Delta R)^2} + \frac{\beta_s}{2(\Delta R)^2}\right\} - \psi_s\left(0, j+1\right)\left\{\frac{1}{\Delta\tau} + \frac{\beta_s}{(\Delta R)^2}\right\} =$$

$$-\psi_s\left(1, j\right)\left\{\frac{\beta_s}{2(\Delta R)^2} + \frac{\beta_s}{2(\Delta R)^2}\right\} + \psi_s\left(0, j\right)\left\{\frac{\beta_s}{(\Delta R)^2} - \frac{1}{\Delta\tau}\right\} \quad (2.70)$$

In tridiagonal terms, this may be rewritten as

$$B\left(0, j\right)\psi_s\left(1, j+1\right) + C\left(m, j\right)\left(0, j+1\right) = D\left(0, j\right) \quad (2.71)$$

Where B(0, j), C(0, j), and D(0, j) has their significance.

3.0 Results and Discussions

The discrete forms of the governing equations modeling the thermal behaviour of the solar our heater are solved numerically. An existing computer programme in BASIC know as EGGINC reported in Enibe (2003) is adapted to implement the numerical solution of the equations and produce performance results. The programme adapted by Enibe (2003) and modified by Obi (2008) to suit his analysis in radial tubes is a sequential modular programme consisting the main programme and some other subprogrammes.The major subgrogramme of the system is named PCM to implement the numerical solutions of the equations modeling the air heater one - dimensional radial heat conduction in a phase change material (PCM) heat storage system encapsulated in cylindrical pipes.Other sub-programmes read and analyze air heater chamber, absorber plate and fins, solar radiation data and compute the particular amount absorbed at any given time. A number of the subroutines calculated heat transfer coefficients for natural or forced convection under various conditions, while a group of other subroutines calculate the thermophysical properties of air at any required temperature. Subprogrammes are also available to implement standard numerical procedures such as function integration using Simpson's rule and the solution of tridiagonal algebraic equations.

The data for experimental and predicted values are reported in Obi (2008) for 12th,13 and 14th June, 2006. as shown in figs a to c,the ambient temperatures are the initial conditions for each day. The collector temperature starts from the initial condition and rises to its maximum value near solar noon corresponding to temperatures of 60^0C to 70^0C. It begins to decrease thereafter as intensity falls. It is observed that these happened in all the figures. The predicted temperatures are plotted alongside the experimental values in figures a to c. In fig a that is June 12th, 2006, the predicted value is about 76^0C and the experimental value is about 70^0C. They are close to within 6^0C. On June 13, 2006, the predicted value is 76^0C and the experimental value is about 56^0C. They are close to within 20^0C.Lastly,on the June 14, 2006, the predicted value is about 74^0C and the experimental value is about 68^0C. They are close to within 4^0C.The above discussion is at the upper limit as shown in the graphs.

At the lower limit in the discussion below the predicted value is 40^0C and the experimental value is 22^0C. They are close to within 18^0C. That is on June 12, 2006. On the 13th June, 2006, the predicted value is 40^0C and the experimental value is 23^0C. They are close to within 17^0C. Lastly,on the 14th June, 2006, the predicted value is about 38^0C and the experimental value is about 25^0C. They are close to within 13^0C.The system operates over

the ambient temperature range of 18^0C - 38^0C and daily global irradiation range of 4.9 – 20.1MJ/m^2-daily. The cumulative useful efficiency is about 13.3%. The reason for the low efficacy is as a result of gloomy and wet days in June. It is observed that from figs a – c, the predicted cure is close enough to the experimental curve within the lower limits of about 6.5^0C and upper limits of about 13^0C.

The predicted values on the system (the solar air heater) for 3 different days in June at Nsukka, Nigeria (latitude 7^0N) of which the initial condition (the ambient temperature and irradiance) for each day are taken from experimental data. The night time performance was not tested. The daily global irradiation covered the range 4.9-20.1MJ/m^2-day, while the ambient temperature over the period varied within the range 18.5-36.0^0C. The initial conditions for each day are taken from the experimental data.
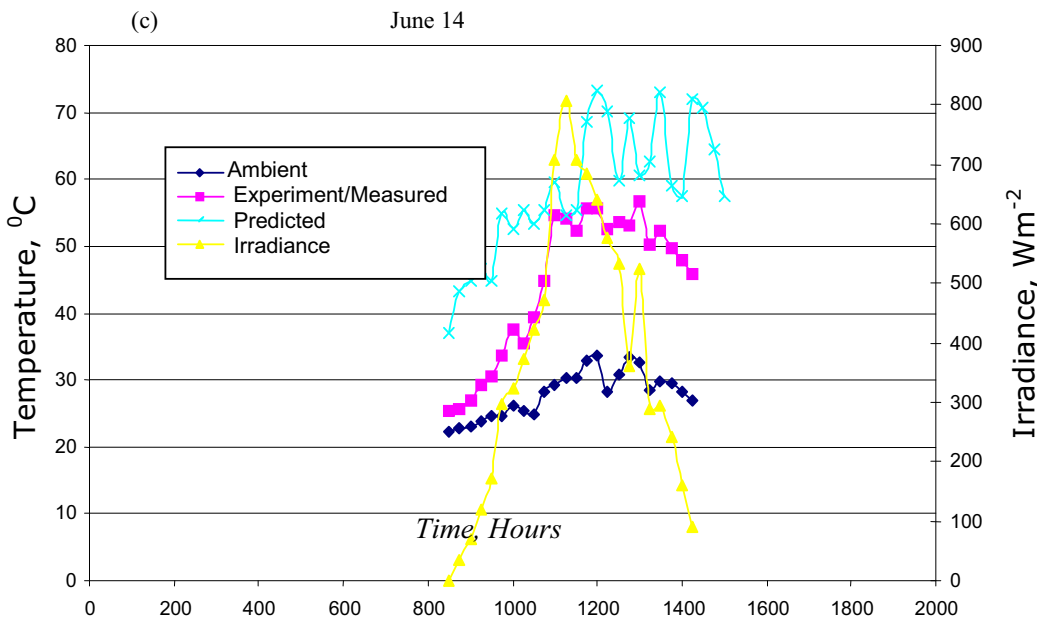
Below are the graphs a – c comparing the experimental temperatures and the predicted temperatures, also shown in the graphs are the irradiance and ambient temperatures.

4.0 Graphs (figs)
Graphs of Temperatures, Irradiance and Ambient Temperatures
[Comparison between the measured and predicted absorber plate temperatures]

(a) June 12

(b) June 13

(c)   June 14
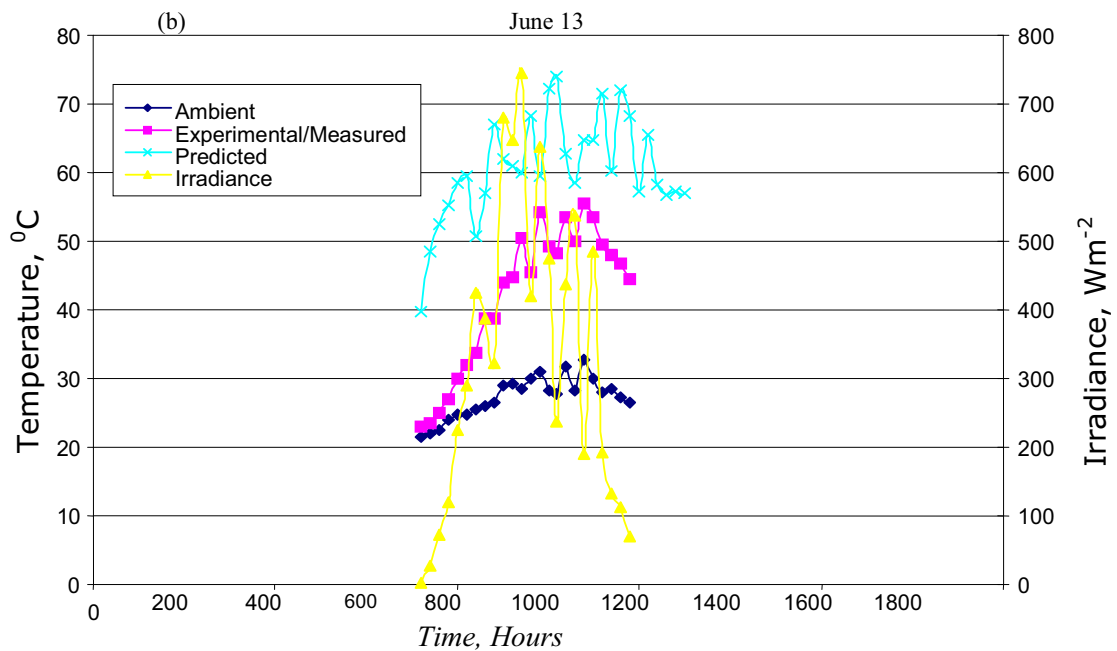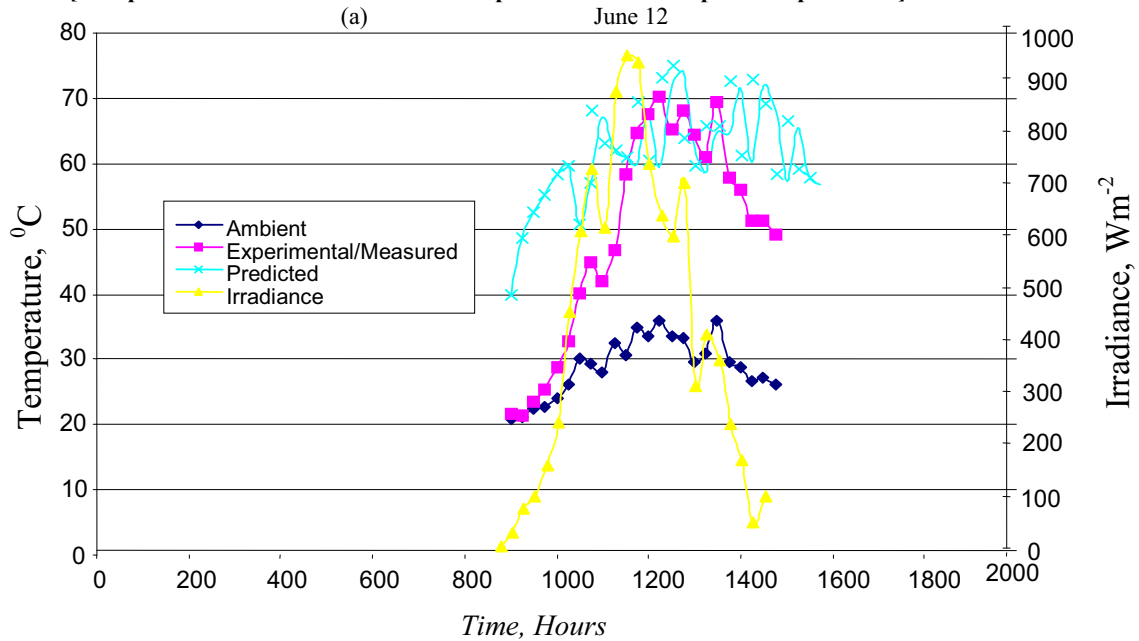
5.0 Conclusion

The finite difference method presented in this work has been employed to deal with one-dimensional heat conduction problems analyzed in cylindrical coordinate. The computer procedure has been found to be reasonable within the trend or within acceptable limits. In this work we consider the phase change material resulting from phase transformations, the melting/freezing problem. Numerical solution of this problem using Crank-Nicolson implicit finite difference scheme is obtained.The one-dimensional radial heat conduction of a phase change materials (PCM) heat storage system encapsulated in cylinders housed in a single glazed flat plate natural circulation solar air heater is presented. Theoretical modeling of the solar air heater predicted the temperature values as compared to the experimental values which agree closely with experimental data to within acceptable limits.

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Nomenclature

A	First term in tridiagonal equation
B	Second term in tridiagonal equation
C	Third term in tridiagonal equation
c_s	Specific heat of the phase change material
D	Tube or pipe diameter
	Last term in tridiagonal equation.
G	Solar radiation intensity, Wm^{-2}
h	Enthalpy, $kJ\ kg^{-1}$
H	Total enthalpy, $kJ\ kg^{-1}$
i	Node identifier in r-direction
j	Node identifier in time direction
k	Thermal conductivity, $Wm^{-1}k^{-1}$
m	Number of nodes in η direction in a pipe containing PCM
r	Radial distance, m
r_o	Centre of the pipe containing the PCM, m
r_m	Radius at the edge or inside radius of the pipe containing the PCM, m
R	Dimensionless distance in the radial direction i.e. r-direction
S	Absorbed radiation, Wm^{-2}
S′	Dimensionless net heat flux
T	Temperature, K
T_f	Fin or plate temperature, K
t	Time, s
t_o	Arbitrary time, s
U	Heat transfer coefficient, $Wm^{-2}K^{-1}$

Greek Letters

α	Thermal diffusivity, $m^2\ s^{-1}$
$β_s$	PCM phase change temperature range factor defined in equations 46 – 50.
δ	Thickness, m
ε	Implicit procedure.
θ	Temperature
	Tangential direction to the pipe containing the PCM (angular), r
ε	Half the difference between upper and lower melting temperature of the PCM, K.
λ	Parameter defined in equation 3-49
λ_1	Parameter defined in equation 4 – 24
λ_2	Parameter defined in equation 4 – 24
ψ	Dimensionless enthalpy
τ	Dimensionless time
Δ	Charge in quantity
φ	Temperature – irradiance function defined in equation.
δ	Partial change or partial derivative

Subscripts

a	Ambient, absorbed
s	Phase change material